



Luminescence Imaging of Solar cells - New Developments

$\mu\Phi$

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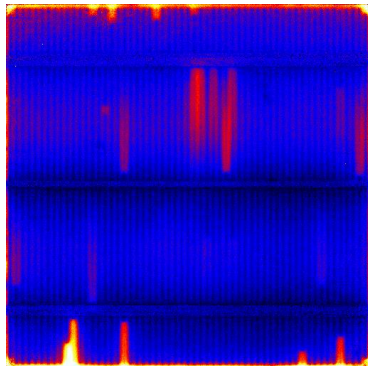
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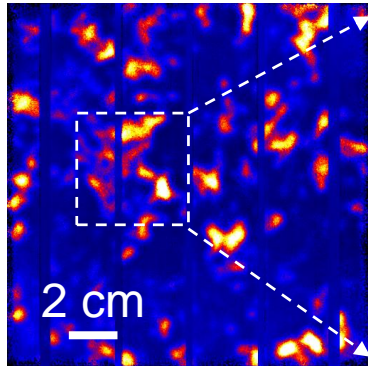
Outline

1. Introduction
2. Why conventional PL- J_{01} imaging is wrong
3. Correct imaging of the calibration constant
4. Easy correction of photon scattering
5. New PL methods for imaging J_{01}
6. Conclusions

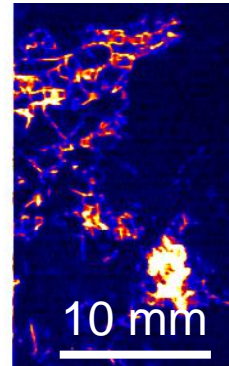
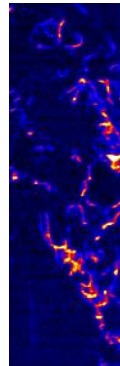
PL- R_s



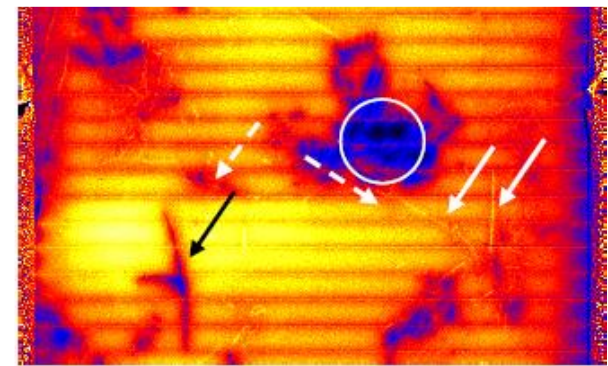
DLIT- J_{01}



PL- J_{01} (enlarged)



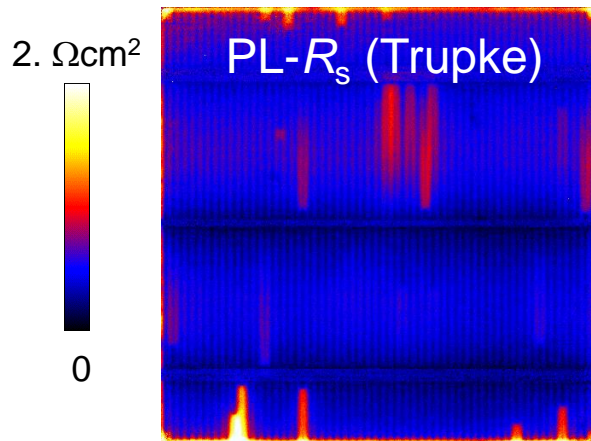
V_{oc} -PL(0.5suns)/ V_{oc} -PL(0.1sun)





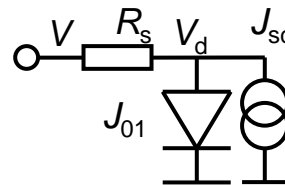
1. Introduction

- Camera-based (Si detector) luminescence imaging (EL + PL) is used for solar cell investigation since 2005^{1,2}
- Starting from 2009, the evaluation was extended to imaging of J_{01} ³⁻⁶
- In 2015 we have shown that this PL-based J_{01} is not correct, since it does not consider the distributed nature of R_s and the action of horizontal balancing currents⁷



$$\Phi = C_i(L_{\text{eff}}) \exp \frac{V_d}{V_T} \quad V_d = V_T \ln \left(\frac{\Phi}{C_i} \right) = V_T (\ln(\Phi) - \ln(C_i))$$

Model of independent diodes (Trupke 2007)



$$V_d = V - R_s \left(J_{01} \exp \frac{V_d}{V_T} - J_{sc} \right)$$

¹T. Fuyuki et al., APL **86** (2005) 262108

²T. Trupke et al., APL **90** (2007) 093506

³M. Glatthaar et al. JAP **105** (2009) 113110

⁴M. Glatthaar et al., PSS RRL **4** (2010) 13

⁵M. Glatthaar et al. JAP **108** (2010) 014501

⁶Chao Shen et al., SOLMAT **109** (2013) 77

⁷O. Breitenstein et al., SOLMAT **137** (2015) 50



1. Introduction

- In 2015 we have found that the usual way for imaging C_i (V_{oc} -PL at 0.1 suns) leads to residual errors in mc cells, an improved method based on linear response principle was proposed.¹
- In 2016 a new method for measuring the PSF for correcting photon scattering in the detector was proposed², enabling accurate Laplacian-based J_{01} imaging.³
- Also in 2016 the „nonlinear Fuyuki“ method was proposed as another alternative PL-based J_{01} imaging method.⁴
- In 2018 it was shown that the luminescence ideality factor may be smaller than unity⁵, and a luminescence-based method to fit a Griddler model to an existing solar cell was proposed.⁶
- This lecture reports about these new developments.

¹O. Breitenstein et al., SOLMAT **142** (2015) 92

²O. Breitenstein et al., J-PV **6** (2016) 522

³F. Frühauf et al., SOLMAT **146** (2016) 87

⁴O. Breitenstein et al., J-PV **6** (2016) 1243

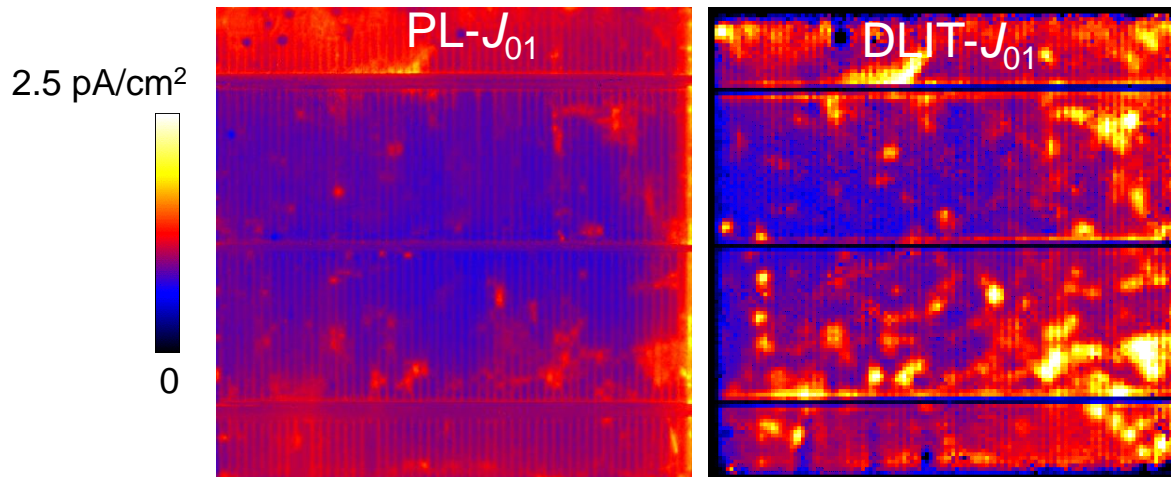
⁵F. Frühauf et al., SOLMAT **180** (2018) 130

⁶F. Frühauf et al., submitted to SOLMAT



2. Why conventional PL- J_{01} imaging is wrong $\mu\Phi$

- It has been found regularly that PL-measured J_{01} images do not agree with DLIT-measured J_{01} images¹
- Chao Shen² has proposed to use n_1 as a global fitting parameter for obtaining a better agreement between PL- and DLIT- J_{01} . However, in our simulations we could not confirm this improvement.

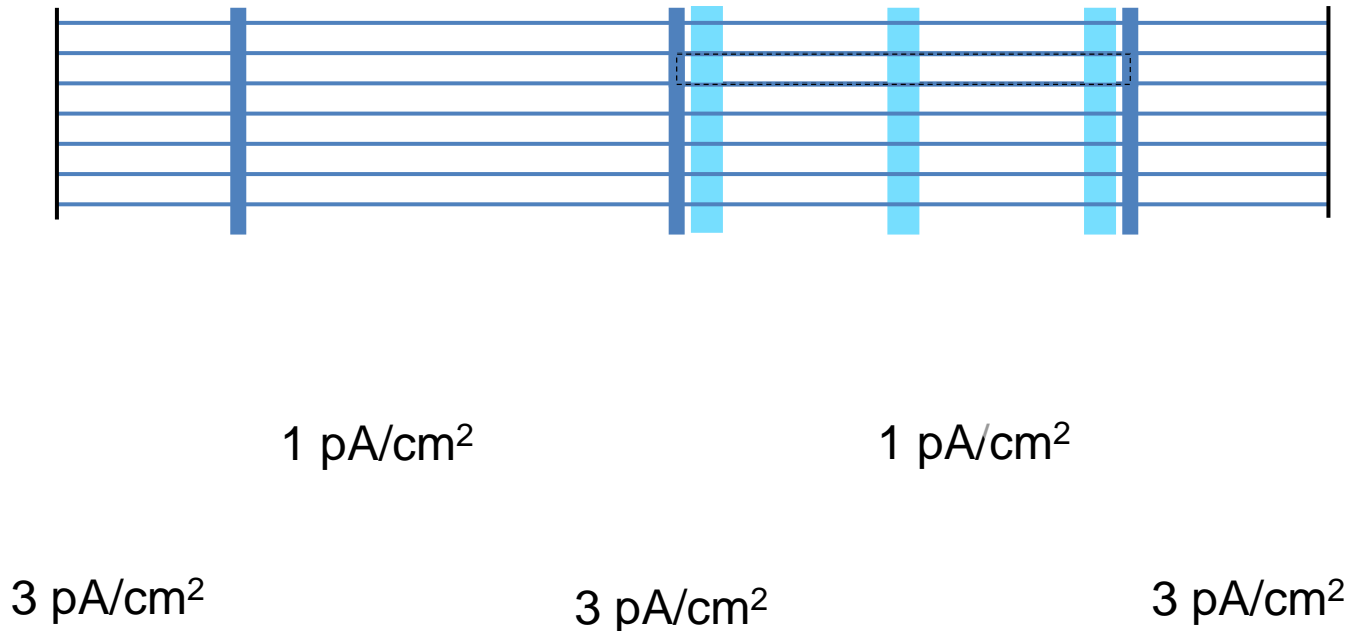


1. O. Breitenstein et al., J-PV **1** (2011) 159
2. Chao Shen et al., SOLMAT **123** (2014) 41



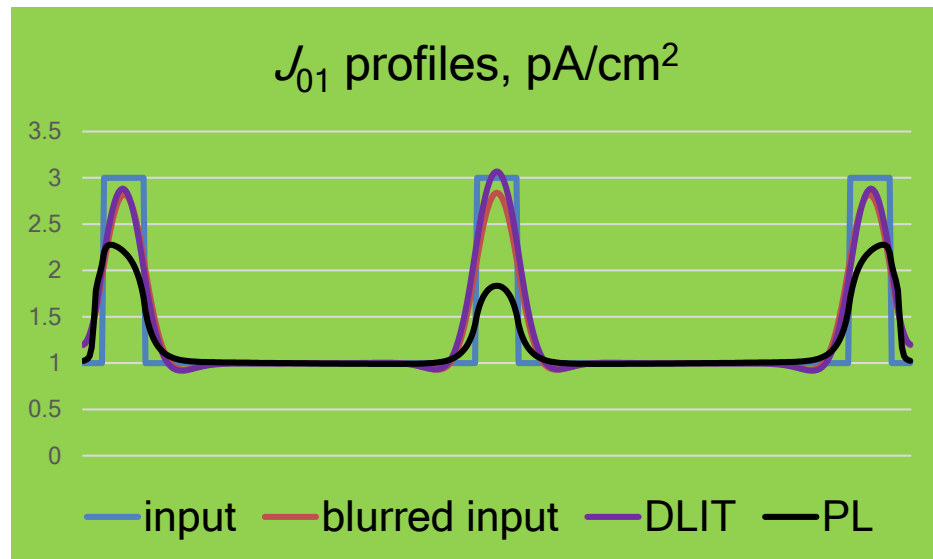
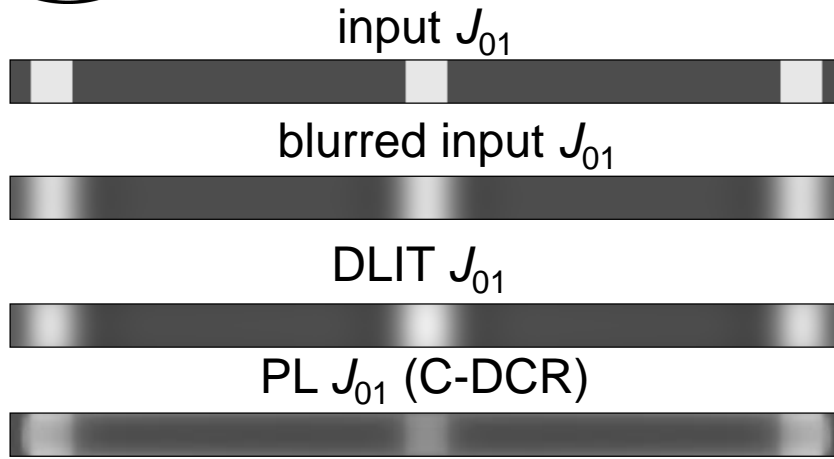
2. Why conventional PL- J_{01} imaging is wrong $\mu\Phi$

- Which of the two results (PL- or DLIT- J_{01}) is correct?
- For answering this question, 2D finite element (SPICE) simulations of a symmetry element of an inhomogeneous solar cell have been performed¹





2. Why conventional PL- J_{01} imaging is wrong $\mu\Phi$

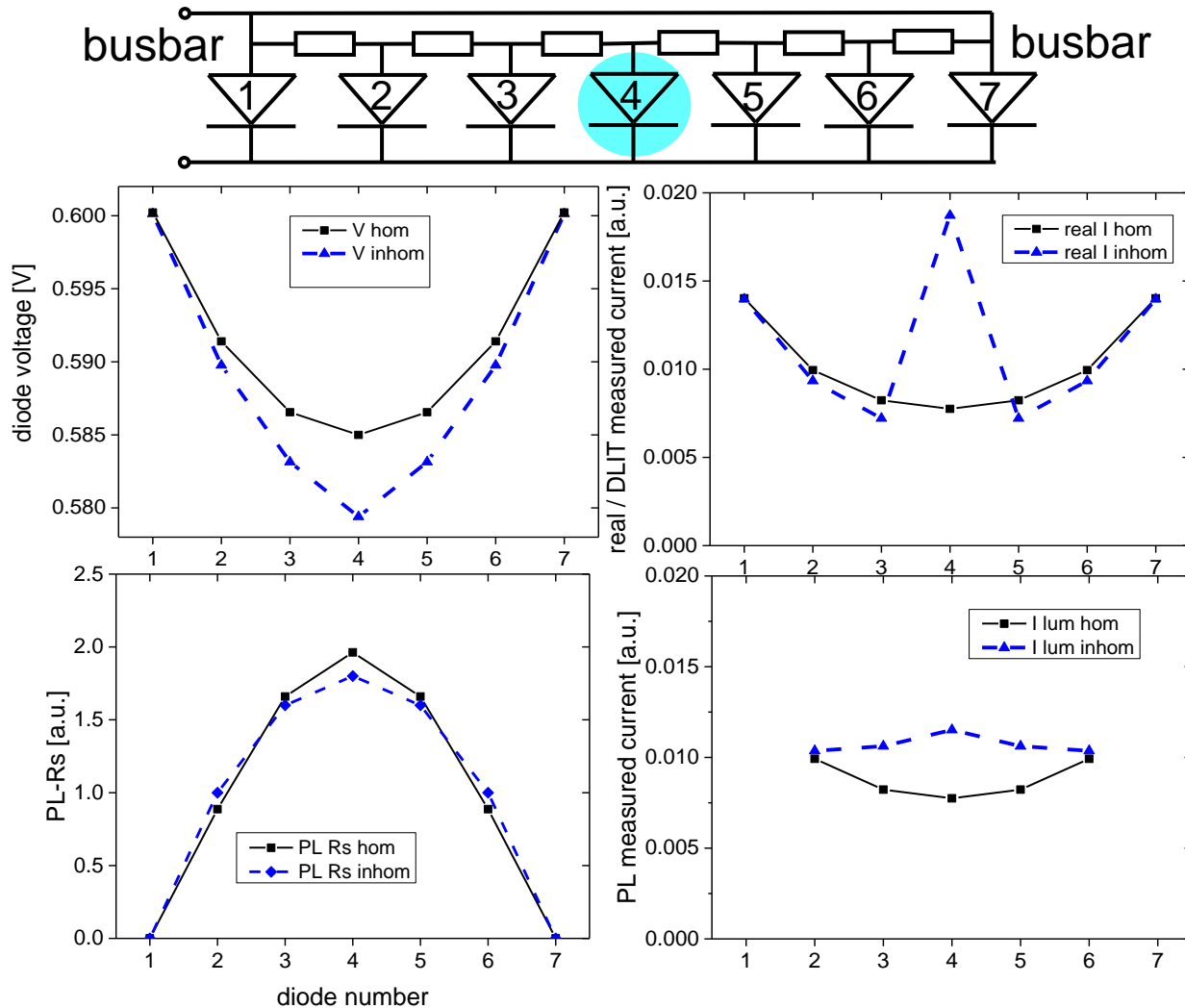


- SPICE simulation of the symmetry element, simulation of PL and DLIT results
- The local maxima of PL- J_{01} calculated by C-DCR appear clearly too weak, they also appear blurred
- This is due to the independent diode model used for C-DCR
- EL/PL can only measure local voltages, the currents follow from the model, which is here too simple
- Also the DLIT evaluation is based on the independent diode model
- However, since in DLIT the current is measured directly, the DLIT results are reliable, except of blurring



2. Why conventional PL- J_{01} imaging is wrong $\mu\Phi$

1-dimensional analog: Resistively coupled diode chain¹

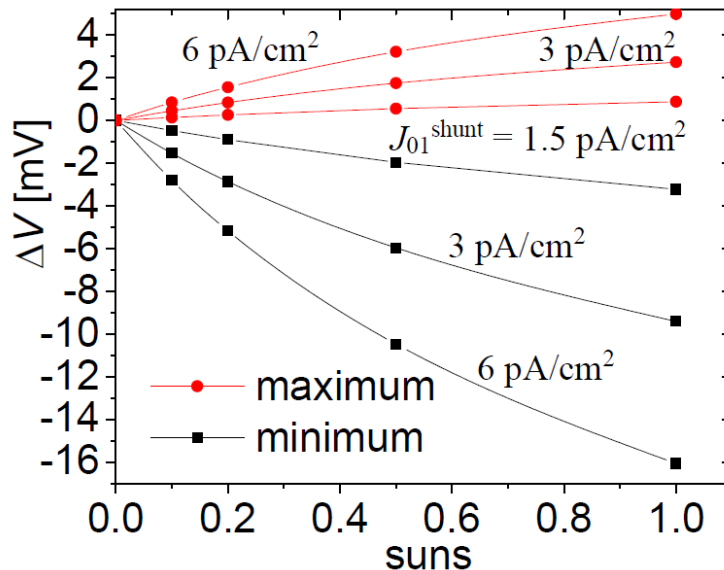


- Only for homogeneous J_{01} , DLIT- and PL-based current imaging results are identical
- If J_{01} shows local maxima, the **resistive intercoupling** leads to horizontal balancing currents, smoothing out the local voltage
- If J is calculated after the usual PL/EL method, **local dark current maxima are underestimated** and the result is blurred



3. Correct imaging of the calibration constant $\mu\Phi$

- SPICE simulation of the symmetry element performed at V_{oc} , various intensities
- Even at V_{oc} (0.1 suns) the local diode voltages are not homogeneously $V_d = V_{oc}$ ¹



$$\Delta V(0.2 \text{ suns}) = \Delta V(0.1 \text{ suns}) * (1 + X)$$

X = nonlinearity parameter, typical value
 $X = 0.86$ for 0.1 and 0.2 suns

- For an unknown cell we do not know $\Delta V(x,y)$
- However, from the linear response principle² we know that this voltage error should be proportional to the illumination intensity $I(\text{suns})$
- For higher intensities the dependence becomes non-linear

¹O. Breitenstein et al., SOLMAT **142** (2015) 92

²J.-M. Wagner et al., Energy Procedia **92** (2016) 255



3. Correct imaging of the calibration constant $\mu\Phi$

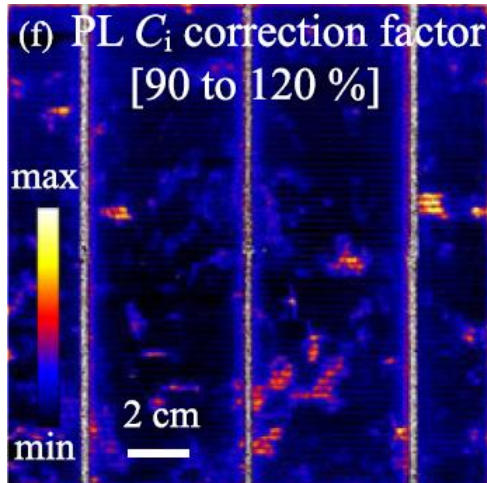
$$\Delta V(0.2 \text{ suns}) = \Delta V(0.1 \text{ suns}) * (1 + X)$$

$$PL^1 = C_i \exp\left(\frac{V_{oc}^1 + \Delta V^1}{V_T}\right)$$

$$PL^2 = C_i \exp\left(\frac{V_{oc}^2 + \Delta V^2}{V_T}\right) = C_i \exp\left(\frac{V_{oc}^2 + (1+X)\Delta V^1}{V_T}\right)$$

$$C_i = \frac{PL^1 \exp\left(\frac{-V_{oc}^1}{V_T}\right)}{\left[\frac{PL^2}{PL^1} \exp\left(\frac{V_{oc}^1 - V_{oc}^2}{V_T}\right)\right]^{\frac{1}{X}}}$$

- This procedure extrapolates C_i to zero illumination intensity, based on the linear response principle¹.
- The only remaining unknown is the nonlinearity parameter X , which may be optimized e.g. by Spice or Griddler simulations³.
- On a usual mc cell, the correction is as large as 20 %, leading to an error of the local $V_{oc}(0.1 \text{ sun})$ of about 5 mV².
- The proposed method provides a clear improvement of the accuracy of C_i imaging. However, **it fails in regions containing ohmic or J_{02} -type shunts** (one-diode model).



¹O. Breitenstein et al., SOLMAT **142** (2015) 92

²O. Breitenstein et al., J-PV **6** (2016) 1243

³F. Frühauf et al., SOLMAT **180** (2018) 130



4. Easy correction of photon scattering $\mu\Phi$

- The importance of **photon scattering in the EL / PL detector** was shown by Walter¹ and the influence of short-pass filtering on the PSF e.g. by Mitchell²
- Due to the limited dynamic range of luminescence detectors, the PSF was measured there by imaging circular apertures of different sizes²
- Teal and Juhl³ have proposed to evaluate the edge spread function (ESF), easily leading to the line spread function (LSF), for obtaining the PSF from one luminescence image. Evaluation method: „backward substitution“
- In cooperation with A. Teal, we have found that this evaluation method leads to certain errors of the PSF and have proposed **an iterative method** for evaluating the LSF⁴
- Our method includes a „correction for diffuse scattering“ and leads to a very exact deconvolution of the input image (zero photon signal in the shadowed region)
- **Our method is meanwhile included in the available „luminescence software suite“⁵**

¹D. Walter et al., Proc. 38th PVSC (2012) 307

²B. Mitchell et al., JAP **112** (2012) 063116

³A. Teal and M. Juhl, Proc. 42nd PVSC (2015)

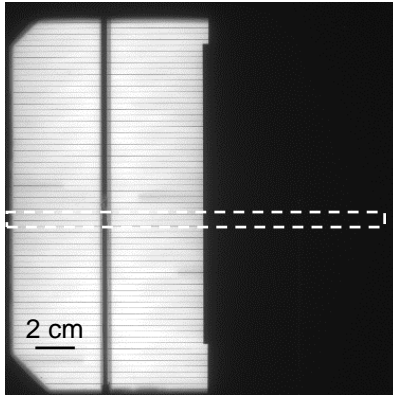
⁴O. Breitenstein et al., J-PV **6** (2016) 522

⁵D.N.R. Payne et al., Comp. Phys. Comm. **215** (2017) 223

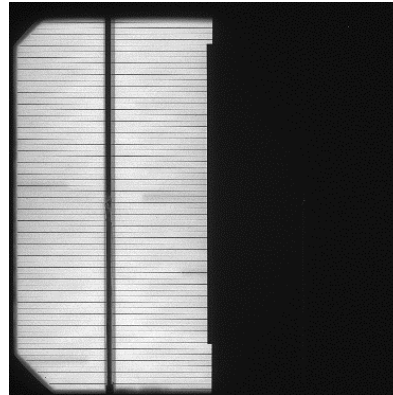


4. Easy correction of photon scattering $\mu\Phi$

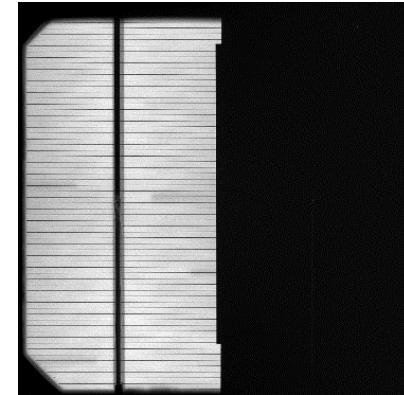
measured EL image
0 to 1 a.u.



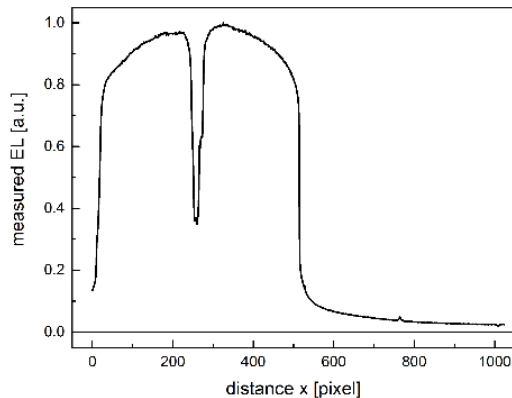
deconvolution
after Teal



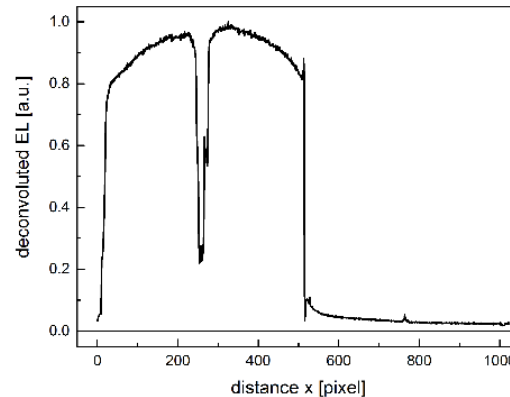
deconvolution after
our method



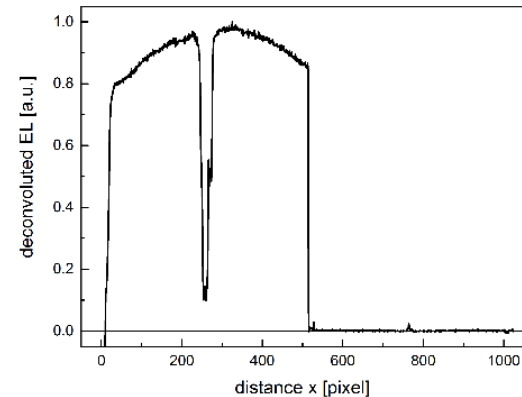
measured EL profile



deconvolution
after Teal



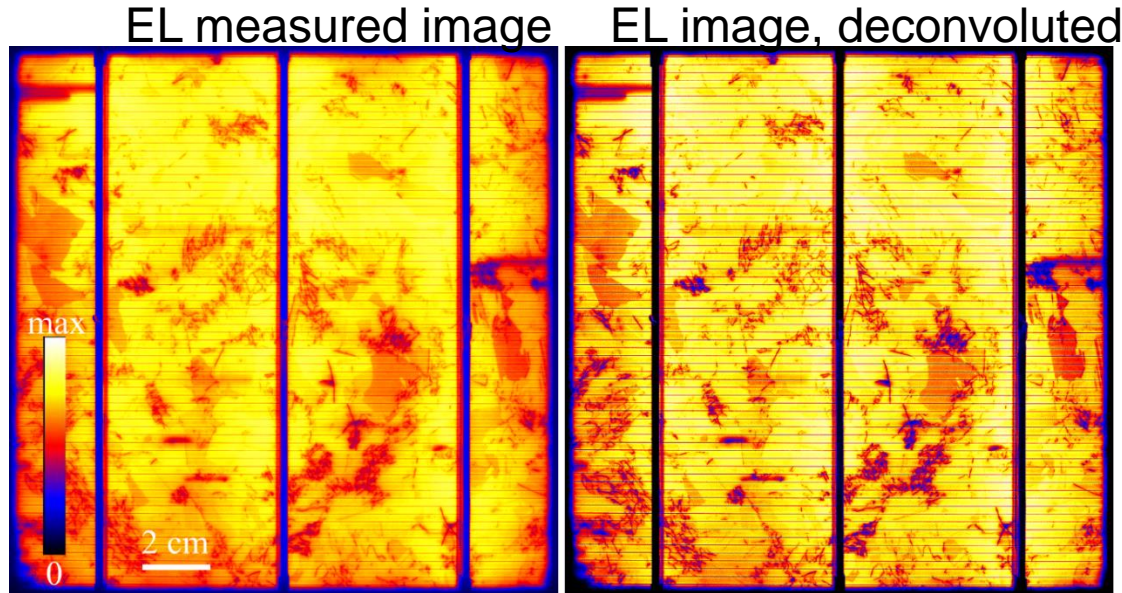
deconvolution after
our method





4. Easy correction of photon scattering $\mu\Phi$

- Effect of deconvolution for a mc standard cell, Si detector without filtering



- If short- or band-pass filtering is used (e.g. 950 to 1000 nm), the effect of light scattering in the detector is strongly reduced, but image acquisition time is increased (x 3 ... 5)
- Then, in many cases, image deconvolution is not necessary anymore.
- If an **InGaAs detector** is used, photon scattering in the detector is negligible, but then **lateral photon scattering in the cell strongly degrades the spatial resolution**¹.

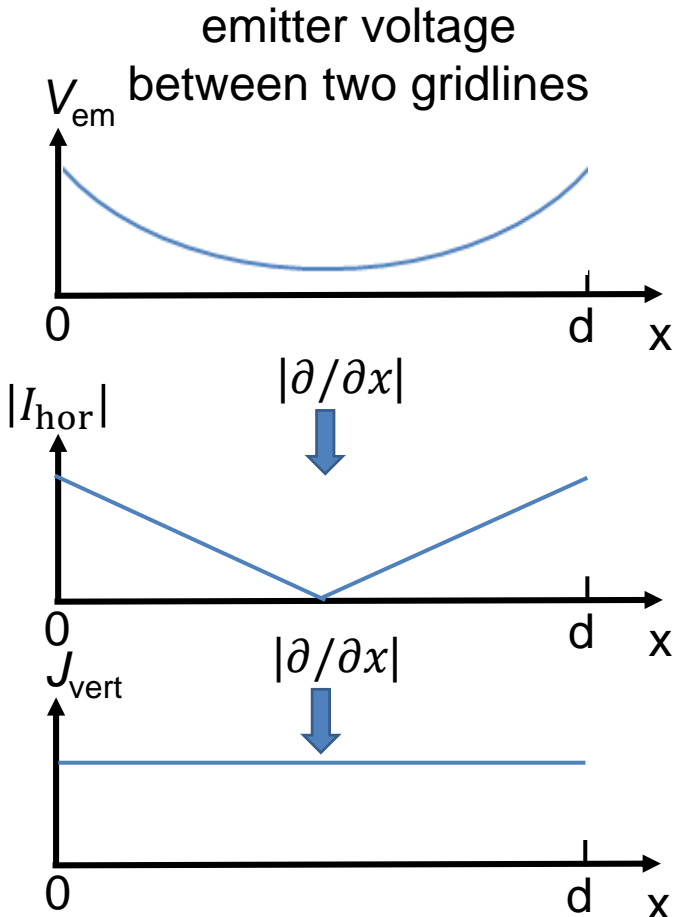


5. New PL methods for imaging $J_{01}\mu\Phi$

1. Laplacian evaluation

Laplacian operator

- First proposed by Glatthaar¹



$$J_{vert}(x, y) = \frac{\partial^2 V_d(x, y)}{Q_{em} \partial x^2} = \frac{\Delta V_d(x, y)}{Q_{em}}$$

$$J_{vert}(x, y) = J_{sc} - J_{01} \exp\left(\frac{V_{oc}^{loc}}{V_T}\right)$$

- $V_{oc}^{loc}(x, y)$ is measured by V_{oc} -PL
- Laplacian evaluation delivers J_{vert} (J_d)
- One-diode model delivers J_{01}
- Main problems: Noise and photon scattering in the detector (blur)



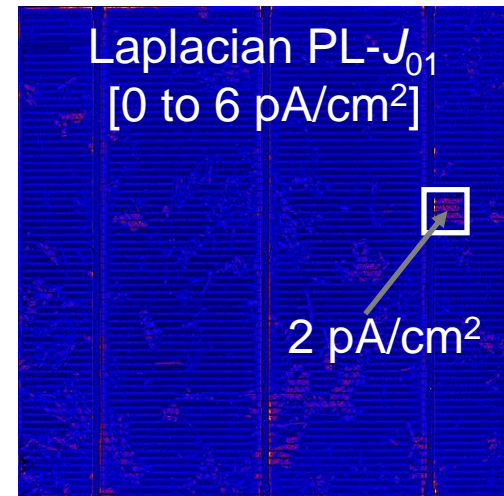
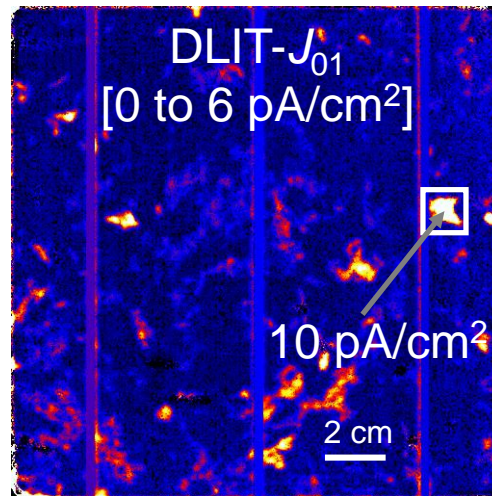
5. New PL methods for imaging $J_{01}\mu\Phi$

1. Laplacian evaluation

- In a pixel image the Laplacian operator [div(grad)] is realized by a pixel sum

$$I_d = (I_{v+} - I_{v-}) + (I_{h+} - I_{h-})$$

- In previous applications of this method data smoothing or pixel binning had to be used. Particularly, J_{01} (or the sheet resistance ρ_s necessary to describe the correct J_{01}) came out a factor of 2...5 too low.

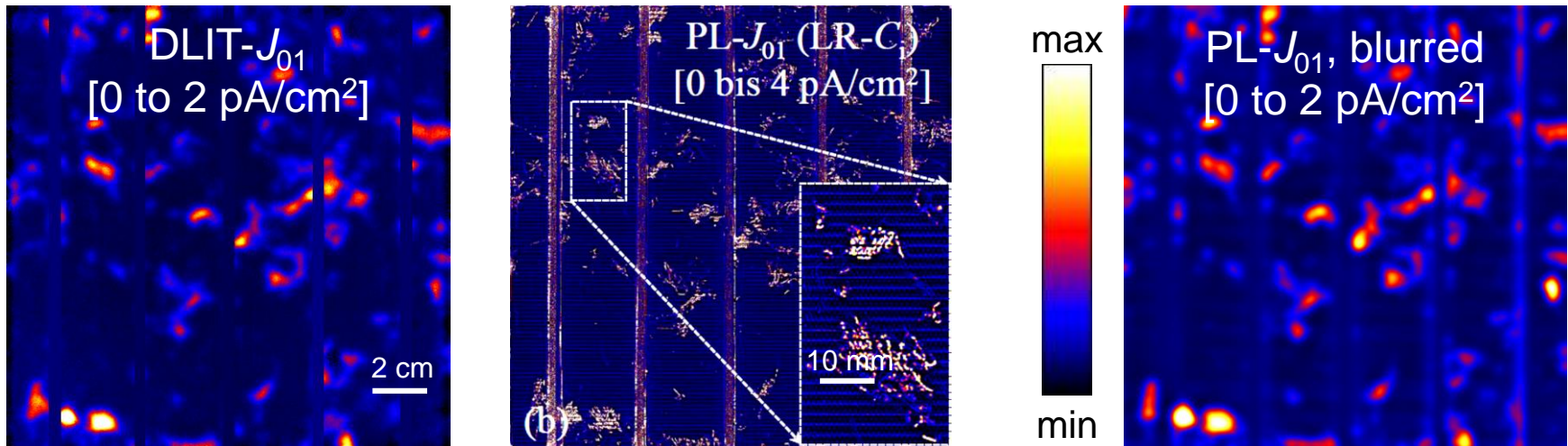




5. New PL methods for imaging $J_{01}\mu\Phi$

1. Laplacian evaluation

- Solution: band-pass filtering + image deconvolution + correct C_i imaging + correcting local diode back voltage¹



- Now Laplacian PL- J_{01} is quantitatively comparable with DLIT- J_{01} , but its spatial resolution is greatly improved.
- Note that this PL evaluation method needs **only one parameter**, which is the emitter sheet resistance ρ_{em} .
- However, **noise is still a problem** for Laplacian PL evaluation.
- Challenge: exclusion of gridlines (spurious signals).



5. New PL methods for imaging $J_{01} \mu \Phi$

2. Nonlinear Fuyuki evaluation

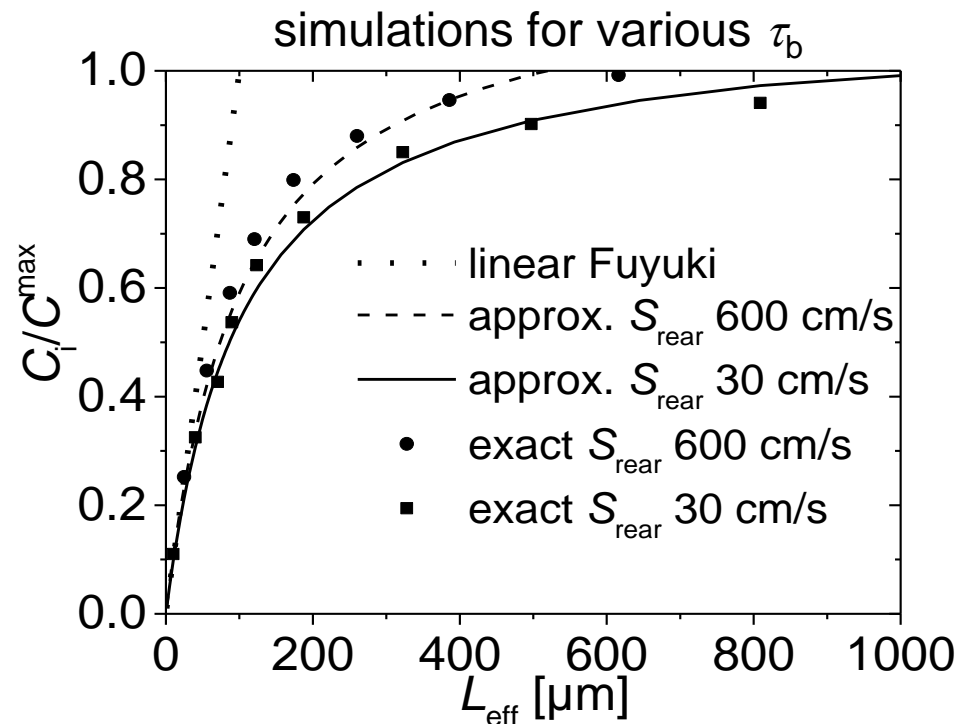
- Fuyuki (APL 2005): „ C_i is proportional to L_d “ (this only holds for very thick cells or low $L_d < 50 \mu\text{m}$).
- Breitenstein¹: non-linear Fuyuki, approximate formula for L_{eff} for low wavelengths

$$\frac{C_i}{C^{\text{max}}} = 1 - \frac{L}{L_{\text{eff}} + L}$$

↑
“information depth”

- C^{max} and L are fitting parameters (fit to DLIT or spectral LBIC).

$$J_{01}^{\text{b}} = \frac{e D n_i^2}{N_A L_{\text{eff}}}$$



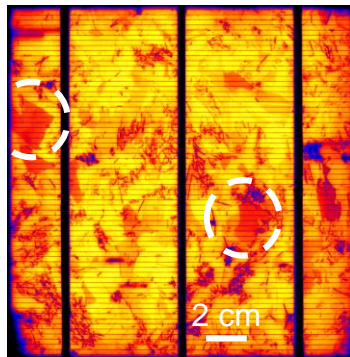


5. New PL methods for imaging $J_{01}\mu\Phi$

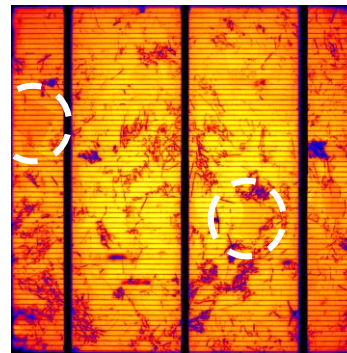
2. Nonlinear Fuyuki evaluation

1. For non-linear Fuyuki evaluation, **band-pass filtering from 950 to 1000 nm** is necessary (inhomogeneous back reflection + theoretical reasons¹).

EL, no filtering

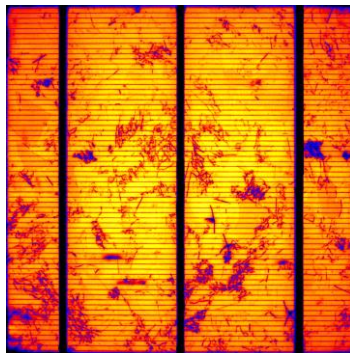


EL, band-pass filtering



2. For non-linear Fuyuki **vignetting correction** (brightness drop at the edges) is necessary.

before correction

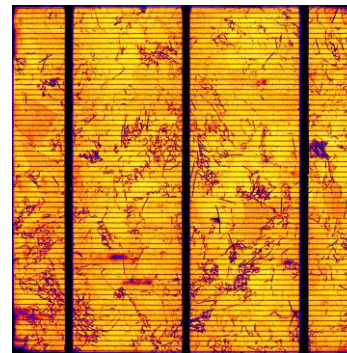


$$C_i^{corr} = \frac{C_i}{\cos^4(A \alpha(x, y))}$$

after correction

↑
imaging
angle

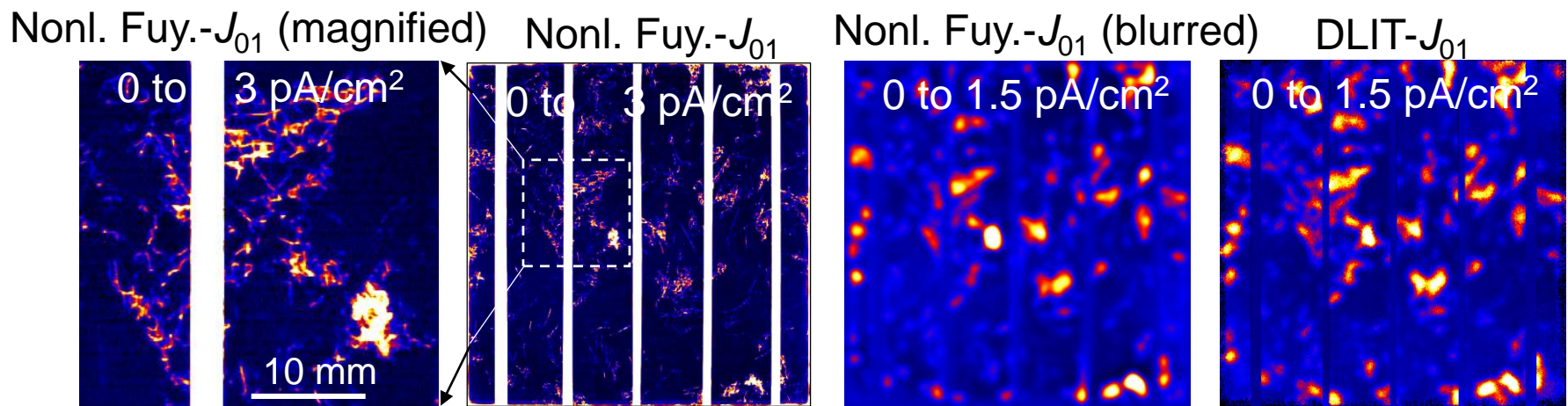
↑
fitting factor (close to 1)



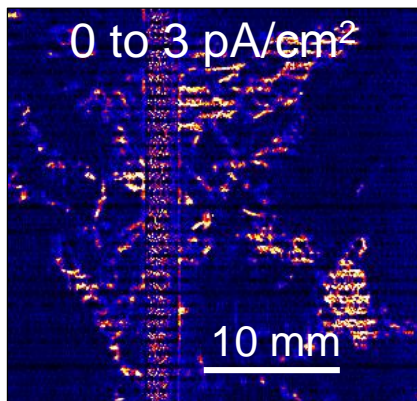


5. New PL methods for imaging $J_{01}\mu\Phi$

2. Nonlinear Fuyuki evaluation



Laplacian PL- J_{01} (magnified)



- If the parameters C_{\max} and L are correctly fitted (e.g. to DLIT- J_{01}), **nonlin. Fuyuki PL- J_{01} images are correct.**
- Their **SNR is clearly better** than that of Laplacian J_{01} .
- Their spatial resolution is also excellent, but shows some residual blurring.
- This blurring is probably due to lateral excess carrier spreading in the cell.



8. Conclusions

- In the last 4 years we have made some significant contributions to PL imaging.
- These are improvements in calculating the luminescence calibration factor C_i and in the calculation of the PSF for correcting photon scattering.
- We have proposed two new methods for high-resolution PL-based imaging of J_{01} / L_{eff} , which may be combined.



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by the German Bundestag

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